

Unitarity in $W_L^+ W_L^-$ scattering without a Higgs particle

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Based on arXiv:1107.1501, with Debmalya Mukhopadhyay

Current status of Higgs searches

Lepton-Photon 2011 announced on Aug 22

"ATLAS and CMS [crudely] taken together exclude all regions for the Standard Model Higgs particle except

- below 145 GeV,
- the range 288-296 GeV, and
- above 464 GeV."

– Matt Strassler, live blogging from Lepton-Photon 2011

What if the Higgs particle really does not exist?

Without a Higgs particle

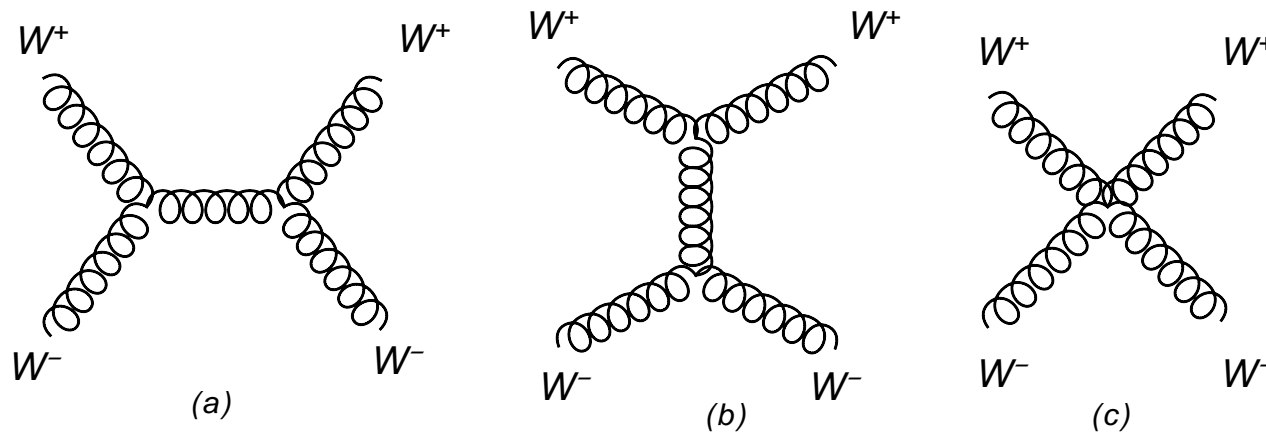
Usual explanations do not work for

- vector boson mass
- $SU(2) \times U(1)$ breaking
 $\rightsquigarrow W - Z$ mass difference
- chiral symmetry breaking
 \rightsquigarrow fermion masses

Fundamental inconsistencies in theory?

$W_L^+ W_L^- \rightarrow W_L^+ W_L^-$ at high energy

SU(2) vector bosons with massive propagators

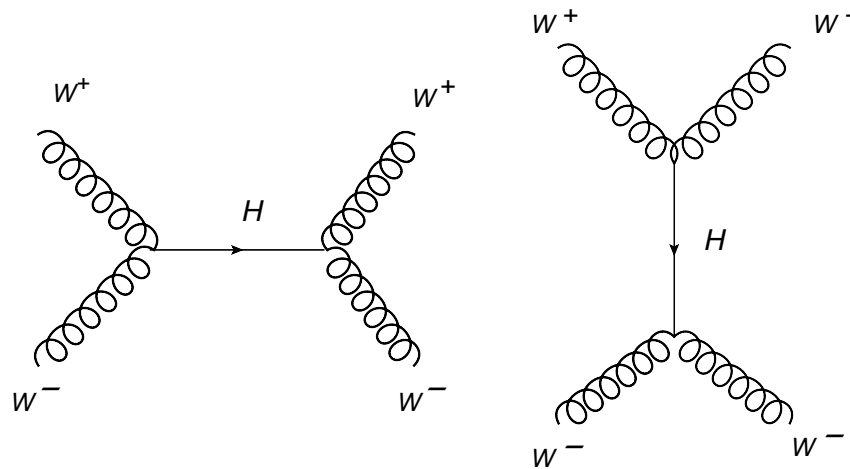


$$\mathcal{M}_W = \mathcal{M}_a + \mathcal{M}_b + \mathcal{M}_c \sim \frac{g^2 P^2}{2m^2} (1 + \cos \theta)$$

Unitarity requires $\text{Re } \mathcal{M} < \frac{1}{2} \Rightarrow$ conflict with QM!

Higgs to the rescue

Two more diagrams if there is a Higgs particle



$$\mathcal{M}_H = \mathcal{M}_s + \mathcal{M}_t \sim -\frac{g^2 P^2}{2m^2}(1 + c)$$

Unitarity is salvaged!

Avoiding the Higgs field

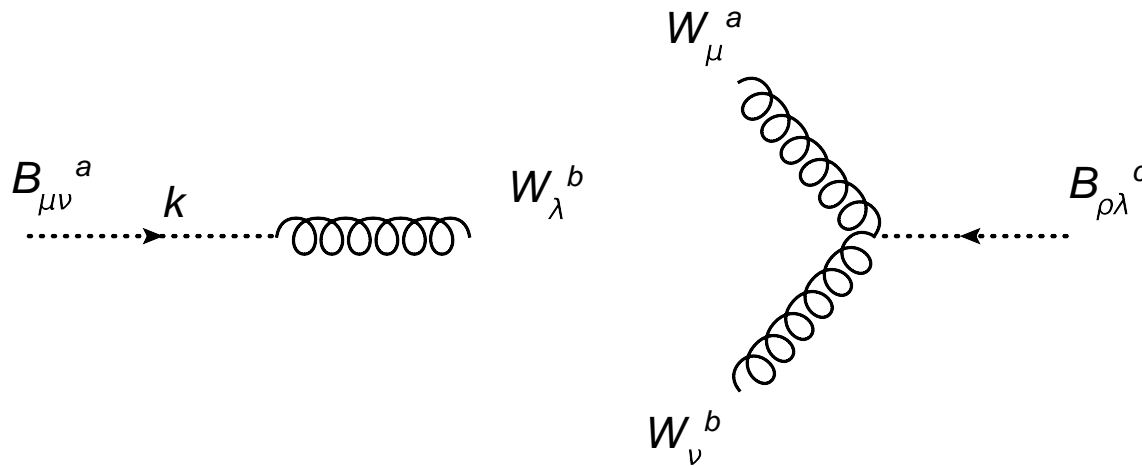
Need two things from any alternative

- Vector boson mass generation
 - Must have propagator $\sim \frac{1}{k^2}$ as $k \rightarrow \infty$
- 3-point vertex with vectors
- (Also other things like renormalizability, absence of ghosts, $SU(2) \times U(1)$ breaking ...)

The Alternative

Triplet of antisymmetric tensor fields $B_{\mu\nu}^a$

Coupled as $m \text{Tr } B \wedge F = \frac{m}{4} \epsilon^{\mu\nu\lambda\rho} B_{\mu\nu}^a F_{\lambda\rho}^a$



$$iV_{\mu\nu,\lambda}^{ab} = -m\delta^{ab}\epsilon_{\mu\nu\lambda\rho}k^\rho$$

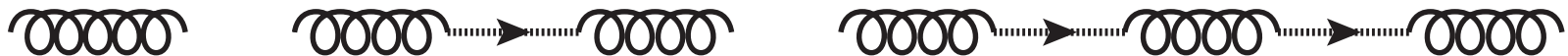
$$iV_{\mu,\nu,\lambda\rho}^{abc} = -igmf^{bca}\epsilon_{\mu\nu\lambda\rho},$$

Vector Boson Mass

Start with 'bare' propagators

$$i\Delta_{\mu\mu'}^{ab} = \frac{-ig_{\mu\mu'}}{k^2 + i\varepsilon} \delta^{ab} \quad i\Delta_{\mu\nu,\mu'\nu'}^{ab} = i \frac{g_{\mu[\mu'} g_{\nu']\nu}}{k^2 + i\varepsilon} \delta^{ab},$$

W propagator = sum over B insertions



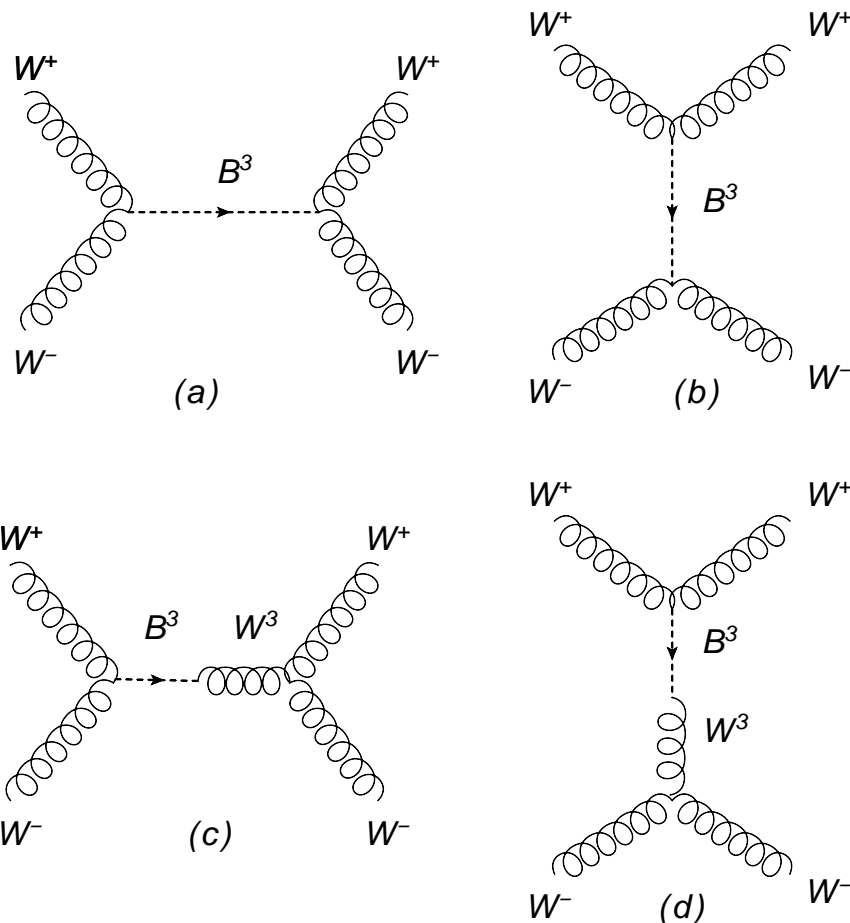
$$\begin{aligned} iD_{\mu\nu} &= i\Delta_{\mu\nu} + i\Delta_{\mu\mu'} iV_{\sigma\rho,\mu'} \frac{i}{4} \Delta_{\sigma\rho,\sigma'\rho'} iV_{\sigma'\rho',\nu'} i\Delta_{\nu'\nu} + \dots \\ &= \frac{-ig_{\mu\nu}}{k^2 + i\varepsilon} \left(1 + \frac{m^2}{k^2} + \frac{m^4}{k^4} + \dots \right) = \frac{-ig_{\mu\nu}}{k^2 - m^2 + i\varepsilon}, \end{aligned}$$

Mass Generation: Comments

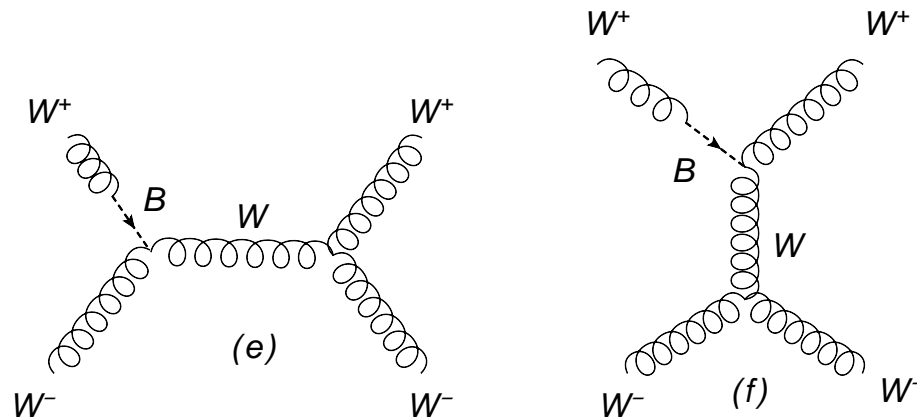
- Massive vector propagator $\sim 1/k^2$ at large k
- B acts like the Goldstone mode of the Higgs field
- B does not take a vev
- All three vectors have the same mass
- $SU(2)$ is unbroken
- No Higgs or Higgs-like particle

B -mediated diagrams

Using the vertices from the $B \wedge F$ term



B-mediated diagrams cont'd



$$\begin{aligned}
 \mathcal{M}_{a+b} + 2\mathcal{M}_{c+d} + 4\mathcal{M}_{e+f} &\sim \left(-\frac{3}{2} + 3 - 2\right) \frac{g^2 P^2}{m^2} (1 + c) \\
 &= -\frac{g^2 P^2}{2m^2} (1 + c) = -\mathcal{M}_W
 \end{aligned}$$

Unitarity is safe! Can we go home now?

Kinetic Term for B

B propagator came from

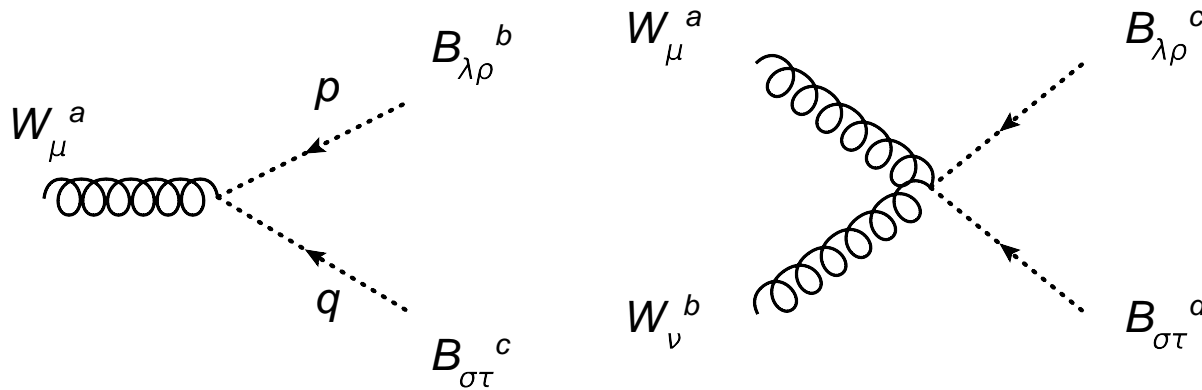
$$\frac{1}{12} H_{\mu\nu\lambda}^a H^{a\mu\nu\lambda}, \quad H_{\mu\nu\lambda} = D_\mu B_{\nu\lambda} + D_\nu B_{\lambda\mu} + D_\lambda B_{\mu\nu}$$

Leads to two more vertex rules

→ from $dB \cdot [A, B]$

→ from $[A, B] \cdot [A, B]$

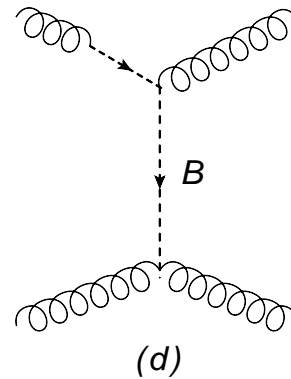
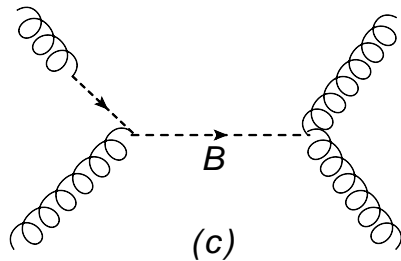
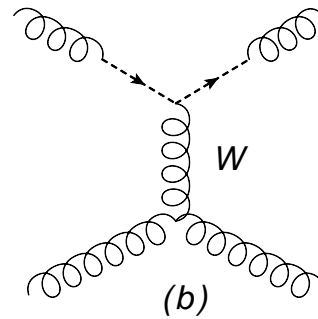
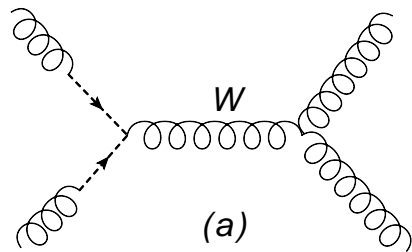
More vertices



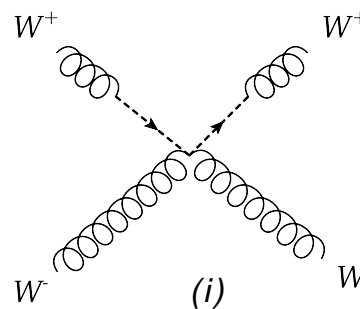
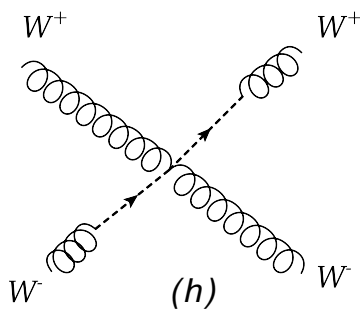
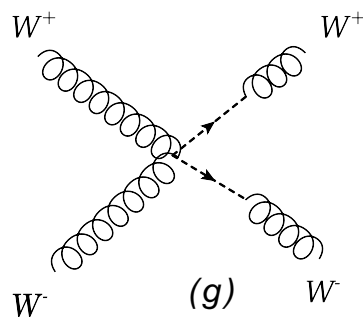
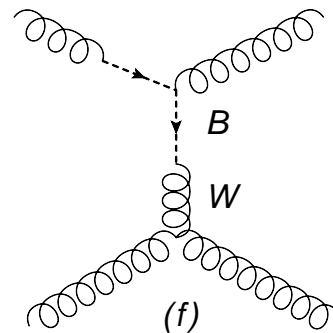
$$\begin{aligned}
 iV_{\mu,\lambda\rho,\sigma\tau}^{abc} &= g f^{abc} \left[(p - q)_\mu g_{\lambda[\sigma} g_{\tau]\rho} + p_{[\sigma} g_{\tau][\lambda} g_{\rho]\mu} - q_{[\lambda} g_{\rho][\sigma} g_{\tau]\mu} \right] \\
 iV_{\mu,\nu,\lambda\rho,\sigma\tau}^{abcd} &= ig^2 \left[f^{ace} f^{bde} \left(g_{\mu\nu} g_{\lambda[\sigma} g_{\tau]\rho} + g_{\mu[\sigma} g_{\tau][\lambda} g_{\rho]\nu} \right) \right. \\
 &\quad \left. + f^{ade} f^{bce} \left(g_{\mu\nu} g_{\lambda[\sigma} g_{\tau]\rho} + g_{\mu[\lambda} g_{\rho][\sigma} g_{\tau]\nu} \right) \right]
 \end{aligned}$$

And more diagrams

With two internal lines of B



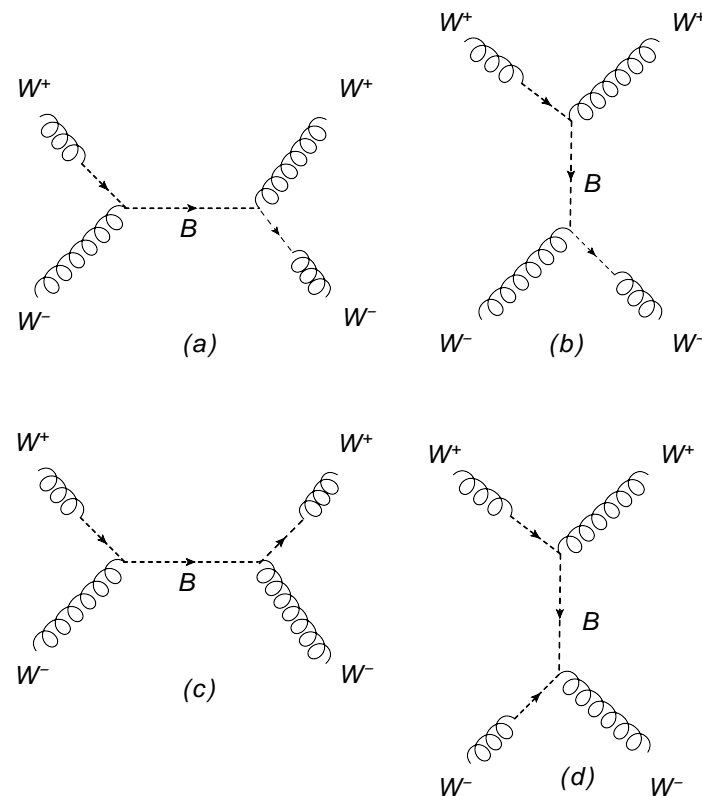
Still with two internal lines of B



$$2\mathcal{M}_{a+b} + 4\mathcal{M}_{c+d} + 4\mathcal{M}_{e+f} + 2\mathcal{M}_{g+h+i} = \frac{4g^2 P^2}{m^2} (1 + 2c + c^2)$$

And more ...

With three internal B lines



$$2\mathcal{M}_a + 2\mathcal{M}_b + 2\mathcal{M}_c + 2\mathcal{M}_d = -\frac{4g^2 P^2}{m^2} (1 + 2c + c^2)$$

Now we are done!

$$\sum_{\text{no Higgs}} \mathcal{M} = \mathcal{O}(P^0)$$

Other issues

- BRST invariance
 - ↪ Using a Stückelberg-type auxiliary field
- Renormalizability
 - ↪ By solving Zinn-Justin Eq.
- $SU(2) \times U(1)$ symmetry breaking
 - ↪ by a soft term (effective term)
- Chiral symmetry breaking
 - ↪ By quantum corrections (finite?)
- Observable effects?
 - ↪ Any process currently involving the Higgs field
 - ↪ Branching ratios should be different



THANK YOU